

Fig. 2 Equivalent circuit of length dz of distributed NLTL.

Following [9], we assume the empirical $v(q_z)$ relation

$$\alpha v(z, t) = \frac{q_z(z, t)}{q_{z0}} + \beta \left(\frac{q_z(z, t)}{q_{z0}} \right)^3 \quad (1)$$

where α , β and q_{z0} characterize HBV. From (1) and Fig. 2 we obtain the nonlinear wave equation

$$\frac{\partial^2}{\partial z^2} [x + \beta x^3] = \frac{1}{u^2} \frac{\partial^2 x}{\partial t^2} \quad (2)$$

where $x(z, t) = q_z(z, t)/q_{z0}$ is the normalized charge/length, $u = 1/(\sqrt{C_{z0} \cdot L_z})$ is a velocity, and $C_{z0} = \alpha q_{z0}$ is the capacitance/length at zero bias.

III. THE NLTL SIMULATION

We simulate the NLTL (a) using a modified FDTD method and (b) by numerical integration.

(a) **The FDTD method** was previously combined with TLM for time-domain simulations of nonlinear devices [10] or transmission lines [11]; the advantages being that the TLM models physical wave propagation, while FDTD offers computational efficiency. Our algorithm resembles [10], except that instead of Maxwell's curl equations, we use the nonlinear wave equation (2) to describe wave propagation. Whereas general FDTD or TLM methods require 6 (field) variables, our algorithm needs only one: the normalized charge x , and is consequently much less computer-intensive. Furthermore, by using an empirical q - v relationship such as (1), it is easy to augment the equivalent circuit for improved accuracy.

Applying FDTD [12] to the wave equation (2), we obtain

$$\begin{aligned} x_{i,n+1} &= 2 \cdot x_{i,n} - x_{i,n-1} \\ &+ k^2 \cdot \mu^2 \cdot \left[\frac{x_{i-1,n-2} - x_{i,n} + x_{i+1,n}}{h^2} \cdot (1 + 3 \cdot \beta \cdot x_{i,n}^2) \right] \\ &+ \frac{3 \cdot k^2 \cdot \mu^2 \cdot \beta \cdot x_{i,n}}{2 \cdot h^2} (x_{i+1,n} - x_{i-1,n})^2. \end{aligned} \quad (3)$$

In (3), the normalized charge $x(z, t)$ is discretized in

space and time as

$$x_{i,n} = x(ih, nk)$$

where $h = \Delta z$ and $k = \Delta t$ are space and time increments, and i and n are integers.

Fig. 3 shows a complete equivalent circuit of the NLTL.

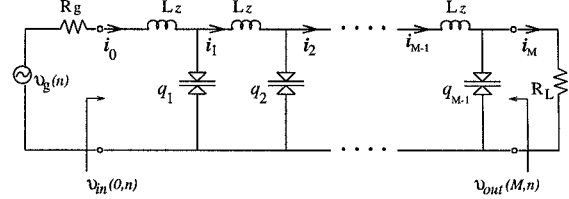


Fig. 3 A complete equivalent circuit of the NLTL.

The boundary conditions are

$$v_{in}(0, n) = v_g(n) - i_0 \cdot R_g \quad (4)$$

$$v_{out}(M, n) = i_M \cdot R_L \quad (5)$$

where R_g and R_L are generator and load resistances, and the equivalent circuit consists of M subsections. The physical length of the NLTL is $l = Mh$.

Application of Kirchhoff's current law and FDTD [12] to Fig. 3 results in

$$\begin{aligned} i_0 &= i_1 + q_{z0} \cdot h \cdot \frac{x_{1,n} - x_{1,n-2}}{2 \cdot k} \\ i_1 &= i_2 + q_{z0} \cdot h \cdot \frac{x_{2,n} - x_{2,n-2}}{2 \cdot k} \\ &\vdots \\ i_M &= i_{M-1} - q_{z0} \cdot h \cdot \frac{x_{M-1,n} - x_{M-1,n-2}}{2 \cdot k} \\ i_{M-1} &= i_{M-2} - q_{z0} \cdot h \cdot \frac{x_{M-2,n} - x_{M-2,n-2}}{2 \cdot k} \\ &\vdots \end{aligned}$$

Combining (4) and (5) with (1), the boundary and initial conditions can be satisfied.

(b) **The numerical integration approach** involves simulating the NLTL in Libra 6.0 [13]. The procedure is:

- (i) Use curve fitting to represent the $C(v)$ relation of (1) by the polynomial

$$C(v) = c_0 + c_1 \cdot v + c_2 \cdot v^2 + \dots + c_m \cdot v^m \quad (6)$$

Here $C(v) = dq_z/dv$ and the coefficients, $c_0, c_1, c_2, \dots, c_m$, depend on the constants, α, β and q_{z0} . The number of the coefficients used here was $m=20$.

- (ii) Use (6) as the nonlinear capacitor model in Libra, and build the equivalent circuit of the NLTL.

- (iii) Do time-domain simulations using the "Transient

Bench” numerical integration procedure of Libra.

Fig. 4 compares the NLTL output voltage waveforms as predicted by the two methods when the input is a 30-GHz 1.0-V sinewave. For the FDTD simulation, the equivalent circuit of Fig. 3 used $M = 25$ subsections. Each subsection corresponds to a discrete HBV of area $27 \times 27 (\mu m)^2$, for an NLTL of width $27 \mu m$, and total length of $25 \times 27 \mu m = 0.67 mm$. The parameters of each section are $\alpha=1.0$, $\beta=3.0$, $L_z h=0.05 nH$, $C_{z0} h=0.1 pF$, $R_s h=1.0 \Omega$ (shunt resistance). Fig. 4 shows good agreement between the FDTD and numerical-integration methods.

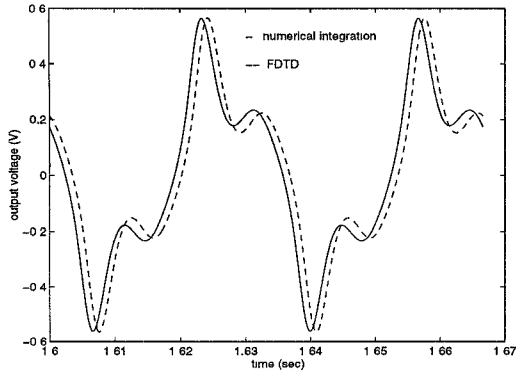


Fig. 4 NLTL time-domain simulation.

IV. RESULTS AND DISCUSSION

Modes of Operation: To investigate the NLTL behavior as a distributed frequency multiplier for different input frequencies f_{in} , we assumed an NLTL of length = 1.35 mm ($M = 50$ subsections), with the parameters per-subsection of $\alpha=1.0$, $\beta=3.0$, $L_z h=0.05 nH$, $C_{z0} h=0.1 pF$, $R_s h=0.5 \Omega$. Fig. 5 shows the simulated output waveform when $f_{in}=29$ GHz.

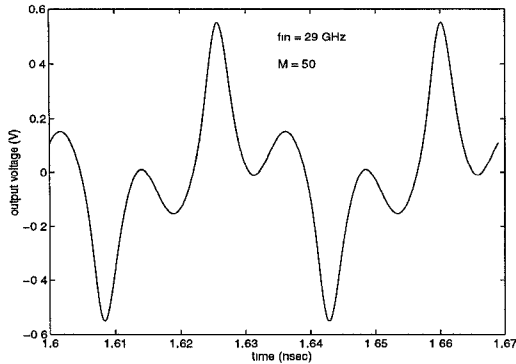


Fig. 5(a) NLTL time-domain response for $l=1.35 mm$, $f_{in}=29 GHz$, and $P_{in}=4 dBm$.

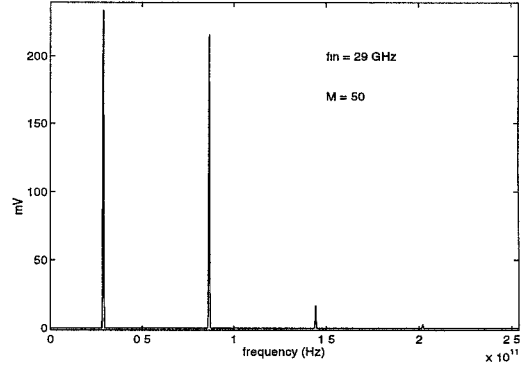


Fig. 5(b) NLTL spectrum for $l=1.35 mm$ ($M=50$), $f_{in}=29 GHz$, and $P_{in}=4 dBm$.

Fig. 5(b) is the corresponding spectrum, showing good tripler operation with $f_{out}=87 GHz$, a 3rd-harmonic / fundamental power ratio of -0.28 dB, and conversion loss of 5.1 dB. All even harmonics are suppressed by the $C(v)$ symmetry, enhancing the conversion efficiency.

Fig. 6(a) depicts the waveform when f_{in} is reduced to 20 GHz: higher-harmonic content is now evident.

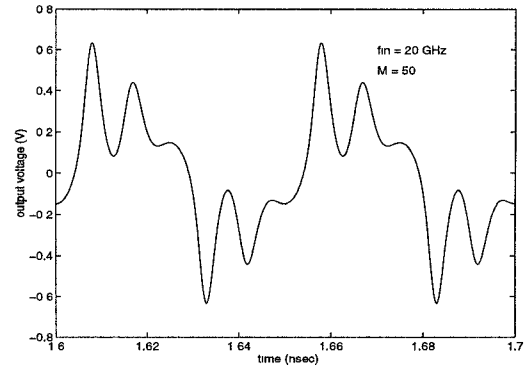


Fig. 6(a) NLTL output waveform for $l=1.35 mm$, $f_{in}=20 GHz$, and $P_{in}=4 dBm$.

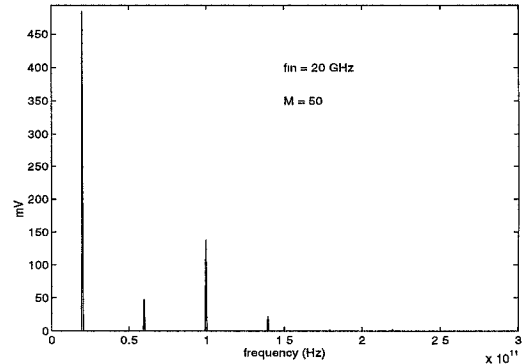


Fig. 6(b) NLTL spectrum for $l=1.35 mm$, $f_{in}=20 GHz$, and $P_{in}=4 dBm$.

As seen in Fig. 6(b), the 5th harmonic is now larger than the 3rd, and the NLTL is operating as a quintupler.

Optimization: To optimize as a tripler, we:

- (i) Set $f_{in} = 30$ GHz and varied the NLTL length over the range $0.54\text{ mm} \leq l \leq 1.89\text{ mm}$ ($20 \leq M \leq 70$)
- (ii) Set $l = 1.35$ mm ($M = 50$) and varied the input frequency over the range $20 \leq f_{in} \leq 30$ GHz.

Fig. 7 shows that for experiment (i) the optimum length is 1.35 mm, yielding a power ratio of -1.69 dB. (compared with the -0.28 dB for $f_{in}=29$ GHz). The power ratio remains better than -3.0 dB over the range $1.08\text{ mm} \leq l \leq 1.62\text{ mm}$ ($40 \leq M \leq 60$).

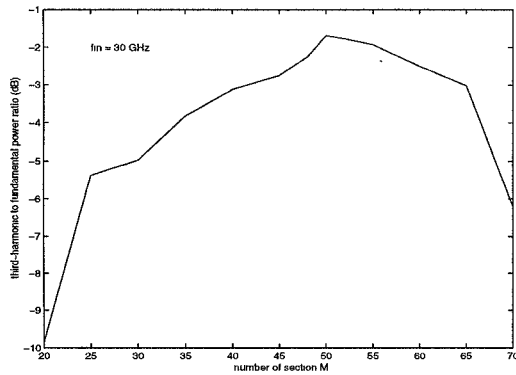


Fig. 7 3rd-harmonic to fundamental power ratio versus NLTL length l (M), when $f_{in}=30\text{GHz}$.

For experiment (ii) Fig. 8 shows an effective input 3-dB bandwidth of 25 to 31.5 GHz, the corresponding output range being 75 to 94.5 GHz.

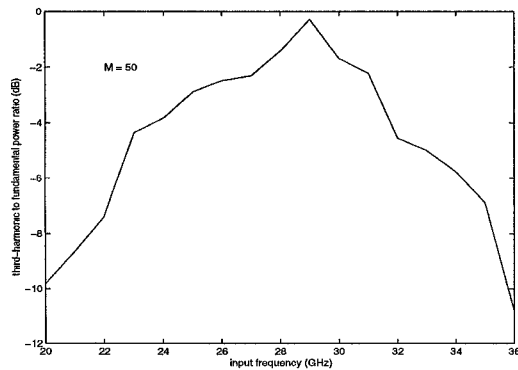


Fig. 8 3rd-harmonic to fundamental power ratio versus input frequency f_{in} , when $l=1.35\text{mm}$ ($M=50$).

V. CONCLUSION

A fully-distributed NLTL incorporating an extended HBV structure is reported. Nonlinear transient analysis is

carried out using both FDTD and numerical-integration methods. The results indicate that this new device is promising as an *untuned* wideband frequency multiplier, with potential in millimeter-wave communications. Further work will be directed to optimizing the parameters of the NLTL and its embedding for efficient harmonic generation for specified output powers levels and bandwidths.

ACKNOWLEDGEMENT

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